

جامعة المنارة

كلية الهندسة

قسم الروبوت و الأنظمة الذكية

مقرر ديناميك الآلات و الاهتزازات الميكانيكية

Planar Kinematics of a Rigid Body
Types of Motion



العام الدراسي 2024-2023

د. محمد خير عبدالله محمد



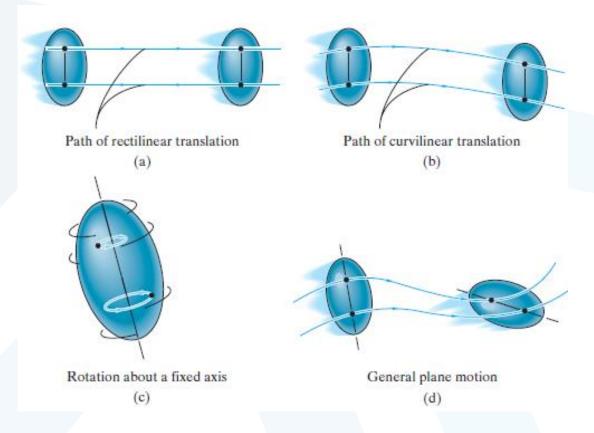
## **Contents**

Planar Rigid-Body Motion
Translation
Rotation about a Fixed Axis
Absolute Motion Analysis



# **Planar Rigid-Body Motion**

The *planar motion* of a body occurs when all the particles of a rigid body move along paths which are equidistant from a fixed plane. There are three types of rigid-body planar motion.



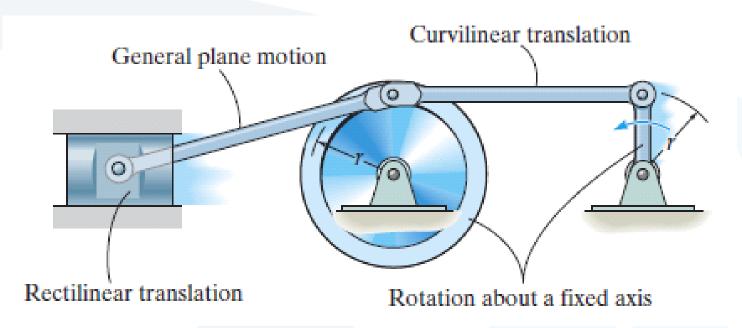


*Translation.* This type of motion occurs when a line in the body remains parallel to its original orientation throughout the motion. When the paths of motion for any two points on the body are parallel lines, the motion is called *rectilinear translation*. If the paths of motion are along curved lines, the motion is called *curvilinear translation*.

Rotation about a fixed axis. When a rigid body rotates about a fixed axis, all the particles of the body, except those which lie on the axis of rotation, move along circular paths.

*General plane motion.* When a body is subjected to general plane motion, it undergoes a combination of translation and rotation. The translation occurs within a reference plane, and the rotation occurs about an axis perpendicular to the reference plane.



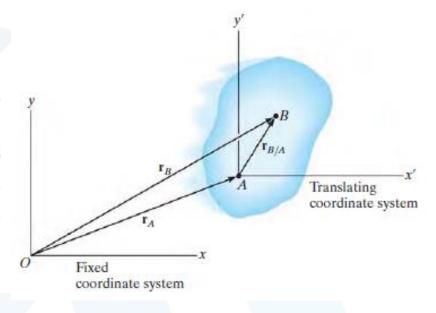




## **Translation**

**Position.** The locations of points A and B on the body are defined with respect to fixed x, y reference frame using *position vectors*  $\mathbf{r}_A$  and  $\mathbf{r}_B$ . The translating x', y' coordinate system is *fixed in the body* and has its origin at A, hereafter referred to as the *base point*. The position of B with respect to A is denoted by the *relative-position vector*  $\mathbf{r}_{B/A}$  (" $\mathbf{r}$  of B with respect to A"). By vector addition,

$$\mathbf{r}_B = \mathbf{r}_A + \mathbf{r}_{B/A}$$





**Velocity.** A relation between the instantaneous velocities of A and B is obtained by taking the time derivative of this equation, which yields  $\mathbf{v}_B = \mathbf{v}_A + d\mathbf{r}_{B/A}/dt$ . Here  $\mathbf{v}_A$  and  $\mathbf{v}_B$  denote absolute velocities since these vectors are measured with respect to the x, y axes. The term  $d\mathbf{r}_{B/A}/dt = \mathbf{0}$ , since the magnitude of  $\mathbf{r}_{B/A}$  is constant by definition of a rigid body, and because the body is translating the direction of  $\mathbf{r}_{B/A}$  is also constant. Therefore,

 $\mathbf{v}_B = \mathbf{v}_A$ 

**Acceleration.** Taking the time derivative of the velocity equation yields a similar relationship between the instantaneous accelerations of *A* and *B*:

$$\mathbf{a}_B = \mathbf{a}_A$$

The above two equations indicate that all points in a rigid body subjected to either rectilinear or curvilinear translation move with the same velocity and acceleration.

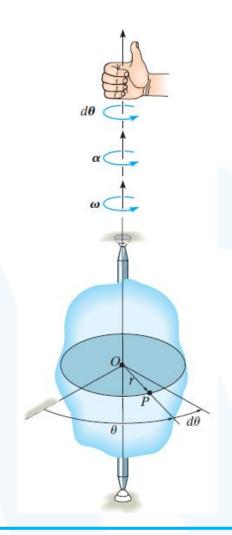


### Rotation about a Fixed Axis

When a body rotates about a fixed axis, any point *P* located in the body travels along a *circular path*. To study this motion it is first necessary to discuss the angular motion of the body about the axis.

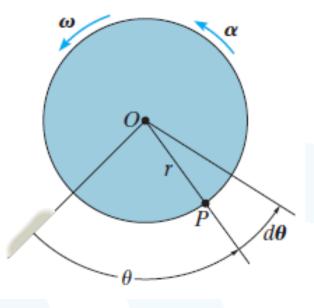
Angular Motion. Since a point is without dimension, it cannot have angular motion. Only lines or bodies undergo angular motion. For example, consider the body shown and the angular motion of a radial line r located within the shaded plane.

Angular Position. At the instant shown, the angular position of r is defined by the angle  $\theta$ , measured from a *fixed* reference line to r.





Angular Displacement. The change in the angular position, which can be measured as a differential  $d\theta$ , is called the *angular displacement*. This vector has a *magnitude* of  $d\theta$ , measured in degrees, radians, or revolutions, where 1 rev =  $2\pi$  rad. Since motion is about a *fixed axis*, the direction of  $d\theta$  is *always* along this axis. Specifically, the *direction* is determined by the right-hand rule; that is, the fingers of the right hand are curled with the sense of rotation, so that in this case the thumb, or  $d\theta$ , points upward. In two dimensions, as shown by the top view of the shaded plane, both  $\theta$  and  $d\theta$  are counterclockwise, and so the thumb points outward from the page.





Angular Velocity. The time rate of change in the angular position is called the *angular velocity*  $\omega$  (omega). Since  $d\theta$  occurs during an instant of time dt, then,

$$(\zeta +) \qquad \qquad \omega = \frac{d\theta}{dt}$$

This vector has a magnitude which is often measured in rad/s. It is expressed here in scalar form since its direction is also along the axis of rotation. When indicating the angular motion in the shaded plane, we can refer to the sense of rotation as clockwise or counterclockwise. Here we have arbitrarily chosen counterclockwise rotations as positive and indicated this by the curl shown in parentheses next to Eq. Realize, however, that the directional sense of  $\omega$  is actually outward from the page.



Angular Acceleration. The angular acceleration  $\alpha$  (alpha) measures the time rate of change of the angular velocity. The magnitude of this vector is

$$\alpha = \frac{d\omega}{dt}$$

$$\alpha = \frac{d^2\theta}{dt^2}$$

The line of action of  $\alpha$  is the same as that for  $\omega$ , however, its sense of direction depends on whether  $\omega$  is increasing or decreasing. If  $\omega$  is decreasing, then  $\alpha$  is called an angular deceleration and therefore has a sense of direction which is opposite to  $\omega$ .

By eliminating dt we obtain a differential relation between the angular acceleration, angular velocity, and angular displacement, namely,

$$(\zeta +)$$

$$\alpha d\theta = \omega d\omega$$



## Constant Angular Acceleration.

$$(\zeta +)$$

$$(\zeta +)$$

$$\omega = \omega_0 + \alpha_c t$$

$$\theta = \theta_0 + \omega_0 t + \frac{1}{2}\alpha_c t^2$$

$$\omega^2 = \omega_0^2 + 2\alpha_c (\theta - \theta_0)$$
Constant Angular Acceleration

Here  $\theta_0$  and  $\omega_0$  are the initial values of the body's angular position and angular velocity, respectively.



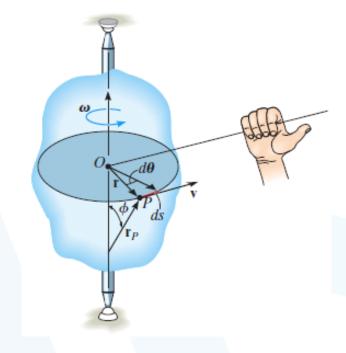
Motion of Point P. As the rigid body rotates, point P travels along a circular path of radius r with center at point O. This path is contained within the shaded plane shown in top view.

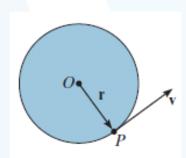
**Position and Displacement.** The position of P is defined by the position vector  $\mathbf{r}$ , which extends from O to P. If the body rotates  $d\theta$  then P will displace  $ds = rd\theta$ .

**Velocity.** The velocity of P has a magnitude which can be found by dividing  $ds = rd\theta$  by dt so that

$$v = \omega r$$

As shown, the direction of v is tangent to the circular path.







Both the magnitude and direction of v can also be accounted for by using the cross product of  $\omega$  and  $r_P$ . Here,  $r_P$  is directed from any point on the axis of rotation to point P. We have

$$\mathbf{v} = \boldsymbol{\omega} \times \mathbf{r}_{P}$$

The order of the vectors in this formulation is important, since the cross product is not commutative, i.e.,  $\boldsymbol{\omega} \times \mathbf{r}_p \neq \mathbf{r}_p \times \boldsymbol{\omega}$ . The fingers of the right hand are curled from  $\boldsymbol{\omega}$  toward  $\mathbf{r}_p$  ( $\boldsymbol{\omega}$  "cross"  $\mathbf{r}_p$ ). The thumb indicates the correct direction of  $\mathbf{v}$ , which is tangent to the path in the direction of motion.

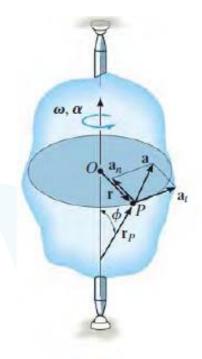


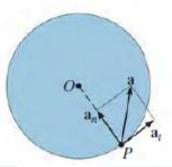
Acceleration. The acceleration of P can be expressed in terms of its normal and tangential components. Applying  $a_t = dv/dt$  and  $a_n = v^2/\rho$ , where  $\rho = r$ ,  $v = \omega r$ , and  $\alpha = d\omega/dt$ , we get

$$a_t = \alpha r$$

$$a_n = \omega^2 r$$

The tangential component of acceleration represents the time rate of change in the velocity's magnitude. If the speed of P is increasing, then  $\mathbf{a}_t$  acts in the same direction as  $\mathbf{v}$ ; if the speed is decreasing,  $\mathbf{a}_t$  acts in the opposite direction of  $\mathbf{v}$ ; and finally, if the speed is constant,  $\mathbf{a}_t$  is zero.







The normal component of acceleration represents the time rate of change in the velocity's direction. The direction of  $a_n$  is always toward O, the center of the circular path

Like the velocity, the acceleration of point P can be expressed in terms of the vector cross product. we have

$$\mathbf{a} = \frac{d\mathbf{v}}{dt} = \frac{d\boldsymbol{\omega}}{dt} \times \mathbf{r}_p + \boldsymbol{\omega} \times \frac{d\mathbf{r}_p}{dt}$$

Recalling that 
$$\alpha = d\omega/dt$$
, and  $(d\mathbf{r}_p/dt = \mathbf{v} = \boldsymbol{\omega} \times \mathbf{r}_p)$  yields 
$$\mathbf{a} = \boldsymbol{\alpha} \times \mathbf{r}_p + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}_p)$$



From the definition of the cross product, the first term on the right has a magnitude  $a_t = \alpha r_p \sin \phi = \alpha r$ , and by the right-hand rule,  $\alpha \times r_p$  is in the direction of  $a_t$ . Likewise, the second term has a magnitude  $a_n = \omega^2 r_p \sin \phi = \omega^2 r$ , and applying the right-hand rule twice, first to determine the result  $\mathbf{v}_p = \boldsymbol{\omega} \times \mathbf{r}_p$  then  $\boldsymbol{\omega} \times \mathbf{v}_p$ , it can be seen that this result is in the same direction as  $a_n$ , shown. Noting that this is also the *same* direction as  $-\mathbf{r}$ , which lies in the plane of motion, we can express  $a_n$  in a much simpler form as  $a_n = -\omega^2 \mathbf{r}$ .

$$\mathbf{a} = \mathbf{a}_t + \mathbf{a}_n$$
$$= \boldsymbol{\alpha} \times \mathbf{r} - \boldsymbol{\omega}^2 \mathbf{r}$$

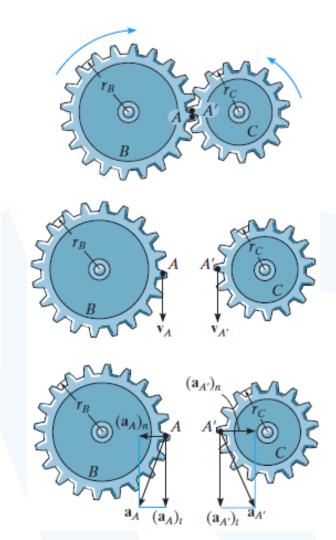
Since  $a_t$  and  $a_n$  are perpendicular to one another, if needed the magnitude of acceleration can be determined from the Pythagorean theorem; namely,  $a = \sqrt{a_n^2 + a_t^2}$ .



If two rotating bodies contact one another, then the *points in contact* move along different circular paths, and the velocity and tangential components of acceleration of the points will be the same: however, the normal components of acceleration will not be the same. For example, consider the two meshed gears. Point A is located on gear B and a coincident point A' is located on gear C. Due to the rotational motion,  $\mathbf{v}_A = \mathbf{v}_{A'}$ , and as a result,  $\omega_B r_B = \omega_C r_C$  or  $\omega_B = \omega_C (r_C/r_B)$ .

Also,  $(\mathbf{a}_A)_t = (\mathbf{a}_{A'})_t$ , so that  $\alpha_B = \alpha_C (r_C/r_B)$ ; however, since both points

Also,  $(\mathbf{a}_A)_t = (\mathbf{a}_{A'})_t$ , so that  $\alpha_B = \alpha_C(r_C/r_B)$ ; however, since both points follow different circular paths,  $(\mathbf{a}_A)_n \neq (\mathbf{a}_{A'})_n$  and therefore, as shown,  $\mathbf{a}_A \neq \mathbf{a}_{A'}$ .



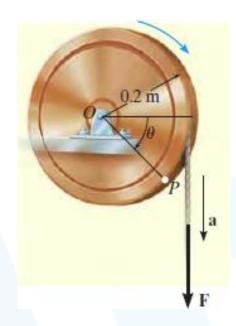


#### EXAMPLE

A cord is wrapped around a wheel, which is initially at rest when  $\theta = 0$ . If a force is applied to the cord and gives it an acceleration  $a = (4t) \text{ m/s}^2$ , where t is in seconds, determine, as a function of time, (a) the angular velocity of the wheel, and (b) the angular position of line OP in radians.

#### SOLUTION

Part (a). The wheel is subjected to rotation about a fixed axis passing through point O. Thus, point P on the wheel has motion about a circular path, and the acceleration of this point has both tangential and normal components. The tangential component is  $(a_p)_t = (4t) \text{ m/s}^2$ , since the cord is wrapped around the wheel and moves tangent to it. Hence the angular acceleration of the wheel is





Using this result, the wheel's angular velocity  $\omega$  can now be determined from  $\alpha = d\omega/dt$ , since this equation relates  $\alpha$ , t, and  $\omega$ . Integrating, with the initial condition that  $\omega = 0$  when t = 0, yields

$$\alpha = \frac{d\omega}{dt} = (20t) \text{ rad/s}^2$$

$$\int_0^{\omega} d\omega = \int_0^t 20t \, dt$$

$$\omega = 10t^2 \text{ rad/s}$$
 Ans.



Part (b). Using this result, the angular position  $\theta$  of OP can be found from  $\omega = d\theta/dt$ , since this equation relates  $\theta$ ,  $\omega$ , and t. Integrating, with the initial condition  $\theta = 0$  when t = 0, we have

$$\frac{d\theta}{dt} = \omega = (10t^2) \text{ rad/s}$$

$$\int_0^{\theta} d\theta = \int_0^t 10t^2 dt$$

$$\theta = 3.33t^3 \text{ rad}$$
Ans.

**NOTE:** We cannot use the equation of constant angular acceleration, since  $\alpha$  is a function of time.



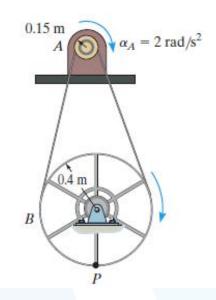
#### EXAMPLE

The motor shown in the photo is used to turn a wheel and attached blower contained within the housing.

If the pulley A connected to the motor begins to rotate from rest with a constant angular acceleration of  $\alpha_A = 2 \text{ rad/s}^2$ ,

determine the magnitudes of the velocity and acceleration of point P on the wheel, after the pulley has turned two revolutions. Assume the transmission belt does not slip on the pulley and wheel.







#### SOLUTION

Angular Motion. First we will convert the two revolutions to radians. Since there are  $2\pi$  rad in one revolution, then

$$\theta_A = 2 \operatorname{rev} \left( \frac{2\pi \operatorname{rad}}{1 \operatorname{rev}} \right) = 12.57 \operatorname{rad}$$

Since  $\alpha_A$  is constant, the angular velocity of pulley A is therefore

$$\omega^2 = \omega_0^2 + 2\alpha_c(\theta - \theta_0)$$

$$\omega_A^2 = 0 + 2(2 \text{ rad/s}^2)(12.57 \text{ rad} - 0)$$

$$\omega_A = 7.090 \text{ rad/s}$$



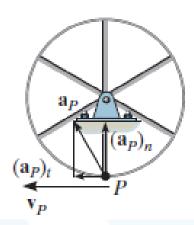
The belt has the same speed and tangential component of acceleration as it passes over the pulley and wheel. Thus,

$$v = \omega_A r_A = \omega_B r_B;$$
 7.090 rad/s (0.15 m) =  $\omega_B$ (0.4 m)  
 $\omega_B = 2.659 \text{ rad/s}$   
 $a_t = \alpha_A r_A = \alpha_B r_B;$  2 rad/s<sup>2</sup> (0.15 m) =  $\alpha_B$ (0.4 m)  
 $\alpha_B = 0.750 \text{ rad/s}^2$ 



### Motion of P. As shown on the kinematic diagram we have

$$v_P = \omega_B r_B = 2.659 \text{ rad/s } (0.4 \text{ m}) = 1.06 \text{ m/s}$$
 Ans.  
 $(a_P)_t = \alpha_B r_B = 0.750 \text{ rad/s}^2 (0.4 \text{ m}) = 0.3 \text{ m/s}^2$   
 $(a_P)_n = \omega_B^2 r_B = (2.659 \text{ rad/s})^2 (0.4 \text{ m}) = 2.827 \text{ m/s}^2$ 



Thus

$$a_P = \sqrt{(0.3 \text{ m/s}^2)^2 + (2.827 \text{ m/s}^2)^2} = 2.84 \text{ m/s}^2$$
 Ans.



# **Absolute Motion Analysis**

A body subjected to general plane motion undergoes a simultaneous translation and rotation. If the body is represented by a thin slab, the slab translates in the plane of the slab and rotates about an axis perpendicular to this plane. The motion can be completely specified by knowing both the angular rotation of a line fixed in the body and the motion of a point on the body. One way to relate these motions is to use a rectilinear position coordinate s to locate the point along its path and an angular position coordinate  $\theta$  to specify the orientation of the line. The two coordinates are then related using the geometry of the problem. By direct application of the time-differential equations v = ds/dt, a = dv/dt,  $\omega = d\theta/dt$ , and  $\alpha = d\omega/dt$ , the motion of the point and the angular motion of the line can then be related.



## **Procedure for Analysis**

The velocity and acceleration of a point P undergoing rectilinear motion can be related to the angular velocity and angular acceleration of a line contained within a body using the following procedure.

#### Position Coordinate Equation.

- Locate point P on the body using a position coordinate s, which is measured from a fixed origin and is directed along the straight-line path of motion of point P.
- Measure from a fixed reference line the angular position  $\theta$  of a line lying in the body.
- From the dimensions of the body, relate s to  $\theta$ ,  $s = f(\theta)$ , using geometry and/or trigonometry.



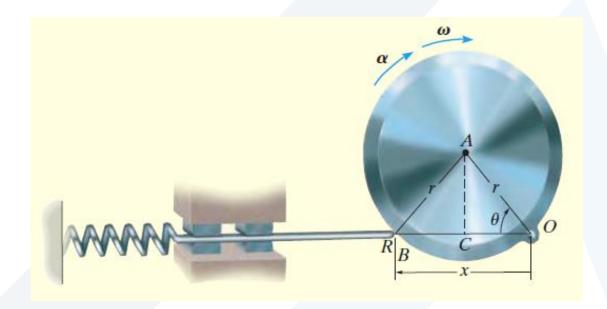
#### Time Derivatives.

- Take the first derivative of  $s = f(\theta)$  with respect to time to get a relation between v and  $\omega$ .
- Take the second time derivative to get a relation between a and  $\alpha$ .
- In each case the chain rule of calculus must be used when taking the time derivatives of the position coordinate equation.



### **EXAMPLE**

The end of rod R maintains contact with the cam by means of a spring. If the cam rotates about an axis passing through point O with an angular acceleration  $\alpha$  and angular velocity  $\omega$ , determine the velocity and acceleration of the rod when the cam is in the arbitrary position  $\omega$ .





#### SOLUTION

**Position Coordinate Equation.** Coordinates  $\theta$  and x are chosen in order to relate the *rotational motion* of the line segment OA on the cam to the *rectilinear translation* of the rod. These coordinates are measured from the *fixed point O* and can be related to each other using trigonometry. Since  $OC = CB = r \cos \theta$ , then

$$x = 2r\cos\theta$$



## Time Derivatives. Using the chain rule of calculus, we have

$$\frac{dx}{dt} = -2r(\sin\theta)\frac{d\theta}{dt}$$

$$v = -2r\omega\sin\theta$$

$$\frac{dv}{dt} = -2r\left(\frac{d\omega}{dt}\right)\sin\theta - 2r\omega(\cos\theta)\frac{d\theta}{dt}$$

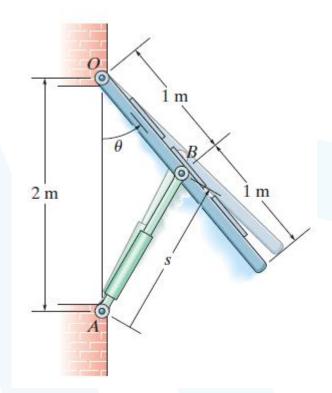
$$a = -2r(\alpha\sin\theta + \omega^2\cos\theta)$$
Ans.

**NOTE:** The negative signs indicate that v and a are opposite to the direction of positive x. This seems reasonable when you visualize the motion.



## EXAMPLE

The large window is opened using a hydraulic cylinder AB. If the cylinder extends at a constant rate of 0.5 m/s, determine the angular velocity and angular acceleration of the window at the instant  $\theta = 30^{\circ}$ .





#### SOLUTION

Position Coordinate Equation. The angular motion of the window can be obtained using the coordinate  $\theta$ , whereas the extension or motion along the hydraulic cylinder is defined using a coordinate s, which measures its length from the fixed point A to the moving point B. These coordinates can be related using the law of cosines, namely,

$$s^{2} = (2 \text{ m})^{2} + (1 \text{ m})^{2} - 2(2 \text{ m})(1 \text{ m})\cos\theta$$
$$s^{2} = 5 - 4\cos\theta \tag{1}$$

When 
$$\theta = 30^{\circ}$$
,

$$s = 1.239 \text{ m}$$



## **Time Derivatives.** Taking the time derivatives of Eq. 1, we have

$$2s\frac{ds}{dt} = 0 - 4(-\sin\theta)\frac{d\theta}{dt}$$
$$s(v_s) = 2(\sin\theta)\omega \tag{2}$$

Since  $v_s = 0.5 \text{ m/s}$ , then at  $\theta = 30^{\circ}$ ,

$$(1.239 \text{ m})(0.5 \text{ m/s}) = 2 \sin 30^{\circ} \omega$$

$$\omega = 0.6197 \text{ rad/s} = 0.620 \text{ rad/s}$$

Ans.



### Taking the time derivative of Eq. 2 yields

$$\frac{ds}{dt}v_s + s\frac{dv_s}{dt} = 2(\cos\theta)\frac{d\theta}{dt}\omega + 2(\sin\theta)\frac{d\omega}{dt}$$
$$v_s^2 + sa_s = 2(\cos\theta)\omega^2 + 2(\sin\theta)\alpha$$

Since  $a_s = dv_s/dt = 0$ , then

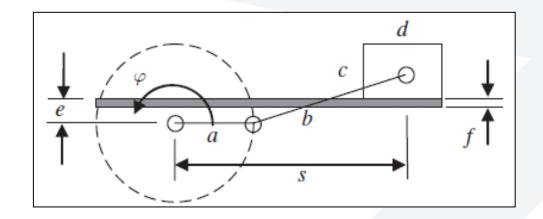
$$(0.5 \text{ m/s})^2 + 0 = 2 \cos 30^{\circ} (0.6197 \text{ rad/s})^2 + 2 \sin 30^{\circ} \alpha$$
  
 $\alpha = -0.415 \text{ rad/s}^2$  Ans.

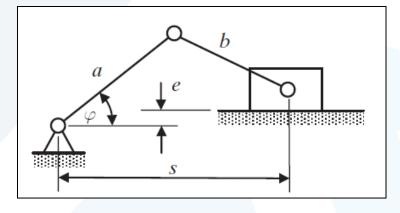
Because the result is negative, it indicates the window has an angular deceleration.



### EXAMPLE

#### Animation of a slider-crank mechanism using Matlab



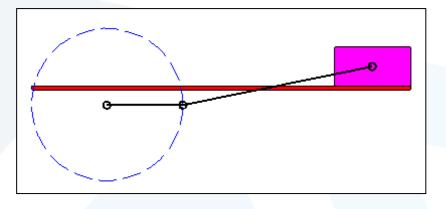


$$s = a\cos(\varphi) + \sqrt{b^2 - (a\sin(\varphi) - e)^2}$$

We shall assume the following numerical values:  $0 \le \varphi \le 2\pi$  at n = 40 equally spaced positions, a = 1, b = 2.5, e = 0.25, c = 0.5, d = 1, and f = 0.06.



```
clc
clear
n = 40;
phi = linspace(0, 2*pi, n);
a = 1;
b = 2.5;
e = 0.25;
nF = 5;
c = 0.5;
d = 1;
f = 0.06;
ax = a*cos(phi);
ay = a*sin(phi);
s = ax + sqrt(b^2-(ay-e).^2);
v = [1.1*min(ax), 1.1*(max(s)+d/2) 1.1*min(ay), 1.1*max(ay)];
xgnd = [min(ax), min(ax), max(s)+d/2, max(s)+d/2, min(ax)];
ygnd = [e-f,e, e, e-f, e-f];
slidery = [e, e+c, e+c, e,e]; % Vertical component of slider is constant
for k = 1:n
   fill(xgnd, ygnd, 'r') % Thin horizontal bar
   hold on
    plot(ax, ay, 'b--', 0, 0, 'ko'); % Dashed circle and center of circle
   sliderx = [s(k)-d/2, s(k)-d/2, s(k)+d/2, s(k)+d/2, s(k)-d/2];
   fill(sliderx, slidery, 'm'); % Slider position
    plot([0 ax(k)], [0 ay(k)], 'ko-', 'LineWidth', 2);
    plot([ax(k), s(k)], [ay(k), e+c/2], 'ko-', 'LineWidth', 2);
   axis(v)
   axis off equal
   SliderCrankFrame(k) = getframe;
   hold off
end
movie(SliderCrankFrame,nF)
```





انتهت المحاضرة