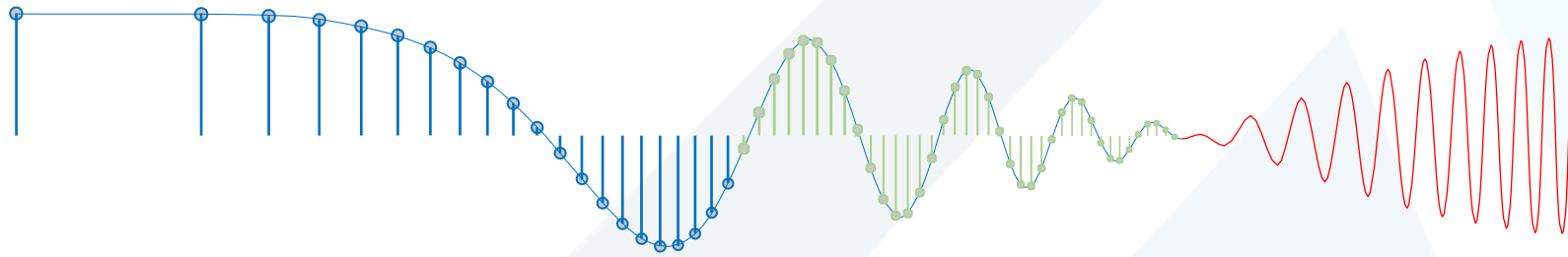


CEDC606: Digital Signal Processing

Lecture Notes 4: Fourier Representation of Signals



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Chapter 4

Fourier representation of signals

1. Sinusoidal signals and their properties
2. Fourier representation of continuous-time signals
3. Fourier representation of discrete-time signals
4. Properties of the discrete-time Fourier transform

Jean-Baptiste Joseph Fourier (1768-1830)



Born 21 March 1768 in Auxerre, Kingdom of France

Died 16 May 1830 (aged 62) in Paris, Kingdom of France

1. Sinusoidal signals and their properties

- Most signals of practical interest can be **described** as a **sum** or **integral** of **sinusoidal** signals.
- However, the exact form of the representation **depends** on whether the signal is **continuous-time** or **discrete-time** and whether it is **periodic** or **aperiodic**.
- For the class of **periodic signals**, such a decomposition is called a **Fourier series**. For the class of **finite energy** signals, the decomposition is called the **Fourier transform**.
- These decompositions are extremely **important** in the **analysis** of **LTI systems** because the **response** of an LTI system to a sinusoidal input signal is a sinusoid of the **same frequency** but of **different** amplitude and phase.

Continuous-time sinusoids

$$x(t) = A\cos(2\pi F_0 t + \phi), \quad -\infty < t < \infty$$

where A is the **amplitude**, ϕ is the **phase** in **radians**, and F_0 is the frequency. The units of F_0 are cycles per second or **Hertz** (Hz). The angular frequency $\Omega_0 = 2\pi F_0$ measured in **radians per second**.

$$A\cos(\Omega_0 t + \phi) = \frac{A}{2} e^{j\phi} e^{j\Omega_0 t} + \frac{A}{2} e^{-j\phi} e^{-j\Omega_0 t}$$

Therefore, we can study the properties of the **sinusoidal signal** by studying the properties of the **complex exponential** $x(t) = e^{j\Omega_0 t}$.

- Let us determine the response $y(t)$ of LTI system to the input $x(t) = e^{j\Omega t}$ using the convolution integral.

$$y(t) = \int_{-\infty}^{\infty} h(\tau)x(t - \tau)d\tau$$

$$y(t) = \int_{-\infty}^{\infty} h(\tau) e^{j\Omega(t-\tau)} d\tau = \int_{-\infty}^{\infty} h(\tau) e^{j\Omega t} e^{-j\Omega\tau} d\tau = \left(\int_{-\infty}^{\infty} h(\tau) e^{-j\Omega\tau} d\tau \right) e^{j\Omega t}$$

$$y(t) = H(j\Omega) e^{j\Omega t}, \quad -\infty < t < \infty$$

- This implies that the **complex exponentials** are **eigenfunctions** of continuous-time LTI systems.
- For a specific value of Ω , the constant $H(j\Omega)$ is an **eigenvalue** associated with the eigenfunction $e^{j\Omega t}$.
- A set of **harmonically** related complex exponential signals, with fundamental frequency $\Omega_0 = 2\pi F_0 = 2\pi/T_0$, is defined by:

$$s_k(t) = e^{jk\Omega_0 t} = e^{j2\pi k F_0 t}, \quad k = 0, \pm 1, \pm 2, \dots$$

- We say that $s_1(t)$ is the **fundamental harmonic** of the set and $s_k(t)$ is the k th **harmonic** of the set. Clearly all harmonics $s_k(t)$ are periodic with period T_0 .

- Furthermore a very important characteristic of harmonically related complex exponentials is the following **orthogonality** property:

$$\int_{T_0} s_m(t) s_n^*(t) dt = \int_{T_0} e^{jm\Omega_0 t} e^{-jn\Omega_0 t} dt = \begin{cases} T_0, & m = n \\ 0, & m \neq n \end{cases}$$

Discrete-time sinusoids

- A discrete-time sinusoidal signal is conveniently obtained by sampling the continuous-time sinusoid at **equally spaced** points $t = nT$:

$$x[n] = x(nT) = A \cos(2\pi F_0 nT + \phi) = A \cos(2\pi \frac{F_0}{F_s} n + \phi)$$

If we define the **normalized frequency** variable: $f = \frac{F}{F_s} = FT$

and the **normalized angular frequency** variable: $\omega = 2\pi f = 2\pi \frac{F}{F_s} = \Omega T$

$$x[n] = A \cos(2\pi f_0 n + \phi) = A \cos(\omega_0 n + \phi), \quad -\infty < n < \infty$$

- The response $y[n]$ of LTI system to the input $x[n] = e^{j\omega n}$:

$$x[n] = e^{j\omega n} \Rightarrow y[n] = H(e^{j\omega})e^{j\omega n}, \quad \text{for all } n$$

which is obtained from z-transform by setting $z = e^{j\omega}$. Thus, the complex exponentials $e^{j\omega n}$ are **eigenfunctions** of discrete-time LTI systems with **eigenvalues** given by the system function $H(z)$ evaluated at $z = e^{j\omega}$.

Periodicity in time: The sequence $x[n] = A\cos(\omega_0 n + \phi)$ is periodic of period N if and only if $f_0 = k/N$, that is, f_0 is a rational number. If k and N are a pair of prime numbers, then N is the fundamental period of $x[n]$. $f_0 = F_0/F_s = k/N = T/T_0$

Periodicity in frequency: The sequence $x[n] = A\cos(\omega_0 n + \phi)$ is periodic in ω_0 with fundamental period 2π and periodic in f_0 with fundamental period one.

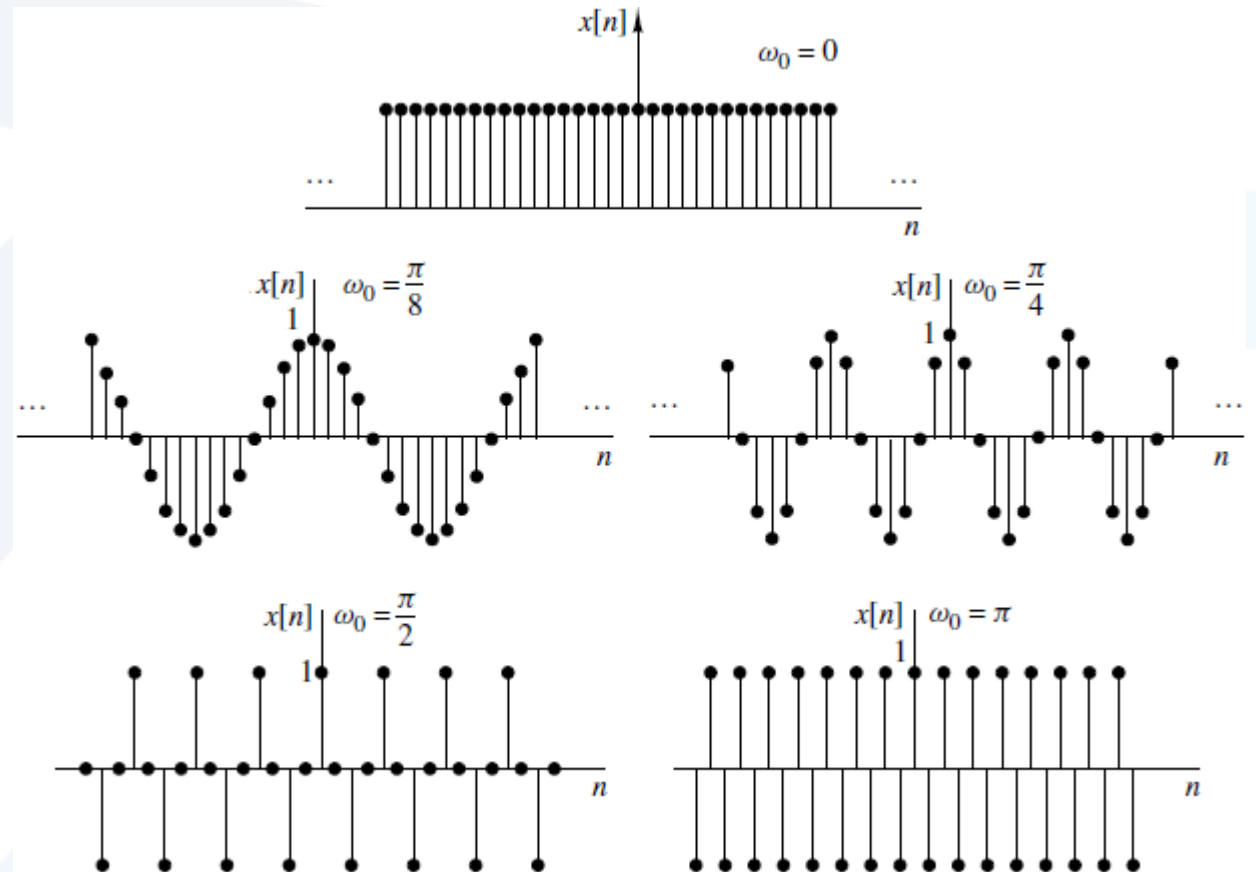
- All distinct sinusoidal sequences have frequencies within an interval of 2π rads. We shall use the **fundamental frequency ranges** $-\pi < \omega \leq \pi$ or $0 \leq \omega < 2\pi$.

- The **rate of oscillation** of a DT sinusoid **increases** as ω_0 **increases** from $\omega_0 = 0$ to $\omega_0 = \pi$. However, as ω_0 increases from $\omega_0 = \pi$ to $\omega_0 = 2\pi$, the oscillations become **slower**.

- Similar properties hold for the DT complex exponentials:

$$s_k[n] = A_k e^{j\omega_k n}, \quad -\infty < n < \infty$$

- For $s_k[n]$ to be **periodic** with fundamental period N , the frequency ω_k should be a rational multiple of 2π , that is, $\omega_k = 2\pi k/N$.



- Therefore, all distinct **complex exponentials** with period N and frequency in the **fundamental range**, have frequencies given by $\omega_k = 2\pi k/N$, $k = 0, 1, \dots, N - 1$. The set of sequences:

$$s_k[n] = e^{j2\pi kn/N}, \quad -\infty < k, n < \infty$$

are periodic both in n (**time**) and k (**frequency**) with fundamental period N .

- There are only N **distinct** harmonically related **complex exponentials** with fundamental frequency $f_0 = 1/N$ and harmonics at frequencies $f_k = k/N$, $0 \leq k \leq N - 1$.

$$s_k[n + N] = s_k[n], \quad (\text{periodic in time})$$

$$s_{k+N}[n] = s_k[n]. \quad (\text{periodic in frequency})$$

$$\sum_{n=\langle N \rangle} s_k[n] s_m^*[n] = \sum_{n=\langle N \rangle} e^{j\frac{2\pi}{N}kn} e^{-j\frac{2\pi}{N}mn} = \begin{cases} N, & k = m \\ 0, & k \neq m \end{cases} \quad (\text{orthogonality property})$$

2. Fourier representation of continuous-time signals

Fourier series for continuous-time periodic signals

- The continuous-time **Fourier series representation (CTFS)** of a **periodic signal**, when it exists, is defined as follows:

1. Synthesis equation:
$$\tilde{x}(t) = \sum_{k=-\infty}^{\infty} c_k e^{jk\Omega_0 t}$$

2. Analysis equation:
$$c_k = \frac{1}{T_0} \int_{T_0} \tilde{x}(t) e^{-jk\Omega_0 t} dt$$

- The set of coefficients $\{c_k\}$ are known as the **Fourier series coefficients**.
- The plot of c_k as a function of frequency $F = kF_0$ (**spectrum**) constitutes a description of the signal in the frequency-domain.

- The plot of $|c_k|$ ($c_k = |c_k|e^{j\phi_k}$) is known as the magnitude spectrum of $\tilde{x}(t)$, while the plot of ϕ_k is known as the phase spectrum of $\tilde{x}(t)$.
- **Parseval's relation:** The **average power** in one period of $\tilde{x}(t)$ can be expressed in terms of the Fourier coefficients using Parseval's relation:

$$P_{av} = \frac{1}{T_0} \int_{T_0} |\tilde{x}(t)|^2 dt = \sum_{k=-\infty}^{\infty} |c_k|^2$$

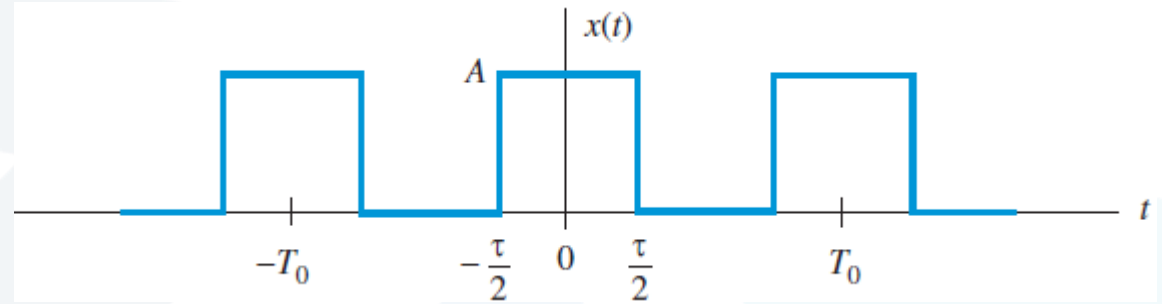
- The value of $|c_k|^2$ provides the **portion** of the average power of signal $\tilde{x}(t)$ that is contributed by the k^{th} **harmonic**.
- The graph of $|c_k|^2$ is known as the **power spectrum** of $\tilde{x}(t)$. (**discrete spectra**).

Fourier transforms for continuous-time aperiodic signals

- We can think of an **aperiodic** signal as a **periodic** signal with **infinite period**.

- We can think of an **aperiodic** signal as a **periodic** signal with **infinite period**. For example let us start with the rectangular pulse train:

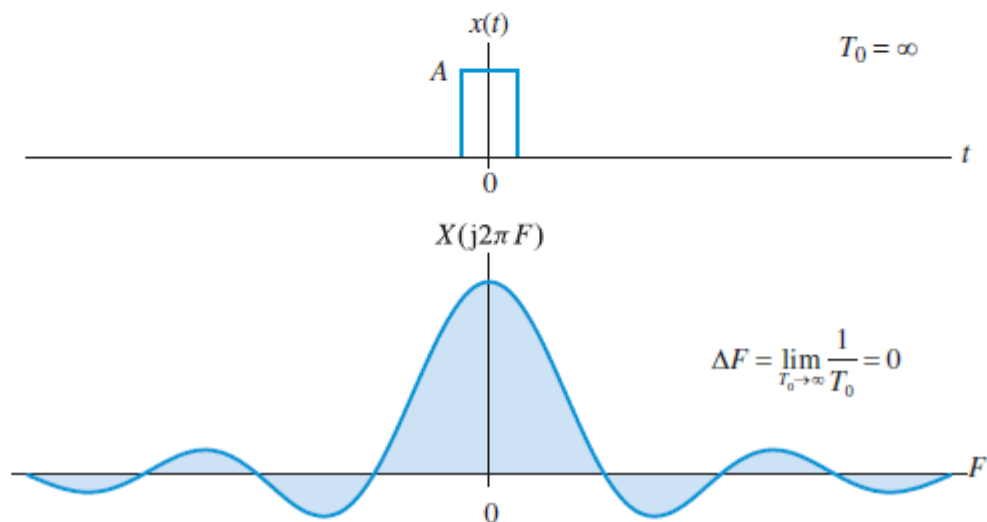
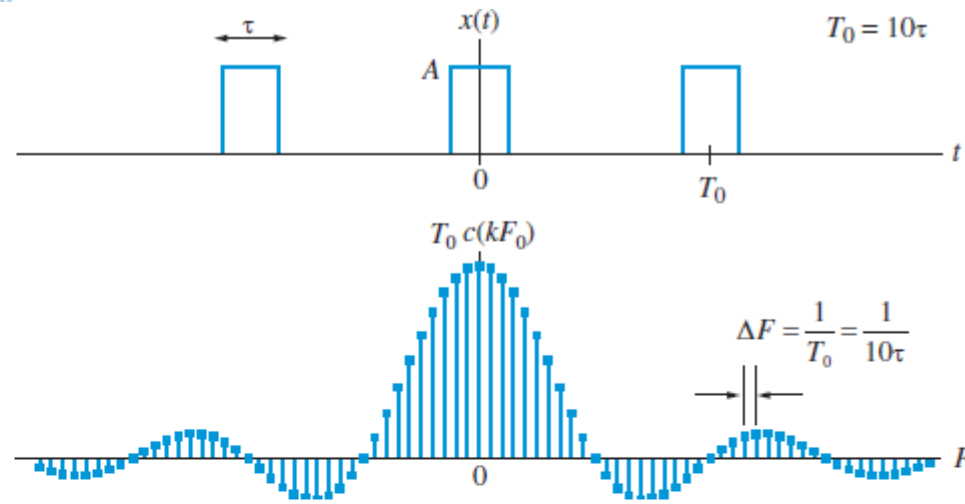
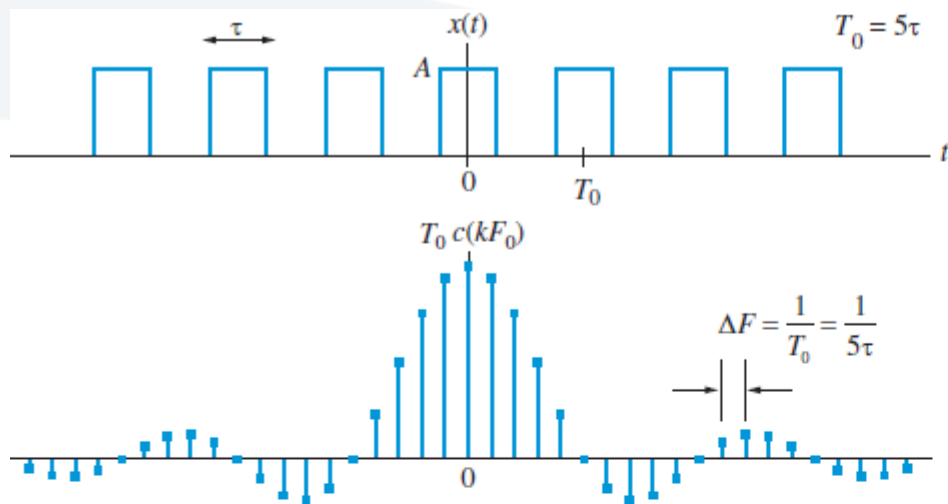
$$x(t) = \begin{cases} A, & |t| < \tau/2 \\ 0, & \tau/2 < |t| < T_0/2 \end{cases}$$



and periodically repeats with period $T_0 > \tau$

$$c_k = \frac{1}{T_0} \int_{-\tau/2}^{\tau/2} A e^{-j2\pi k F_0 t} dt = \frac{A\tau}{T_0} \frac{\sin \pi k F_0 \tau}{\pi k F_0 \tau}, \quad k = 0, \pm 1, \pm 2, \dots$$

- As $T_0 \rightarrow \infty$ (a) in the **time domain** the result is an aperiodic signal corresponding to one period of the rectangular pulse train, and (b) in the **frequency domain** the result is a “continuum” of spectral lines.



- The continuous-time **Fourier transform representation (CTFT)** of an **aperiodic signal**, when it exists, is defined as follows:

1. Synthesis equation: $x(t) = \mathcal{F}^{-1}\{X(j2\pi F)\} = \int_{-\infty}^{\infty} X(j2\pi F) e^{j2\pi Ft} dF$

$$x(t) = \mathcal{F}^{-1}\{X(j\Omega)\} = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\Omega) e^{j\Omega t} d\Omega$$

2. Analysis equation: $X(j2\pi F) = \mathcal{F}\{x(t)\} = \int_{-\infty}^{\infty} x(t) e^{-j2\pi Ft} dt$

$$X(j\Omega) = \mathcal{F}\{x(t)\} = \int_{-\infty}^{\infty} x(t) e^{-j\Omega t} dt$$

- **Parseval's relation:** For **aperiodic** signals with **finite energy**, Parseval's relation is given by:

$$\int_{-\infty}^{\infty} |x(t)|^2 dt = \int_{-\infty}^{\infty} |X(j2\pi F)|^2 dF$$

- The quantity $|X(j2\pi F)|^2$ is known as the **energy-density spectrum** of $x(t)$.

3. Fourier representation of discrete-time signals

Fourier series for discrete-time periodic signals

- The discrete **Fourier series representation (DTFS)** of a **periodic signal**, is defines as follows:

1. Synthesis equation:
$$\tilde{x}[n] = \sum_{k=0}^{N-1} c_k e^{j\frac{2\pi}{N}kn}$$

2. Analysis equation:
$$\tilde{c}_k = \frac{1}{N} \sum_{n=0}^{N-1} \tilde{x}[n] e^{-j\frac{2\pi}{N}kn}$$

- Parseval's relation:** The **average power** in one period of $\tilde{x}[n]$ can be expressed as:

$$P_{av} = \frac{1}{N} \sum_{n=0}^{N-1} |\tilde{x}[n]|^2 = \sum_{k=0}^{N-1} |\tilde{c}_k|^2$$

- The value of $|\tilde{c}_k|^2$ provides the **portion** of the average power of signal $\tilde{x}[n]$ that is contributed by the k^{th} **harmonic**.
- There are only N distinct harmonic components.
- The graph of $|\tilde{c}_k|^2$ as a function of $f = k/N$, $\omega = 2\pi k/N$, or simply k , is known as the **power spectrum** of the periodic signal $\tilde{x}[n]$.

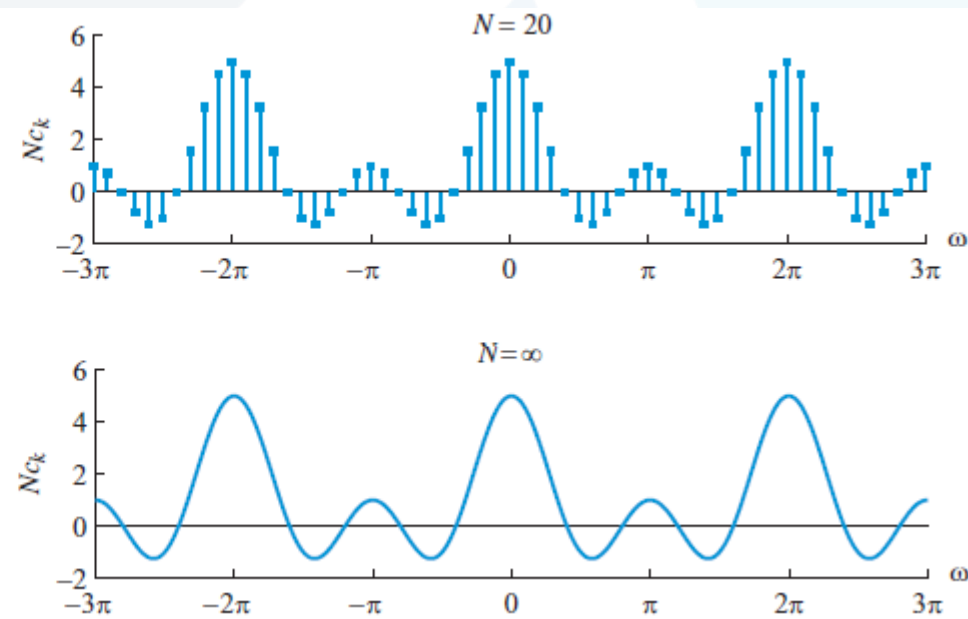
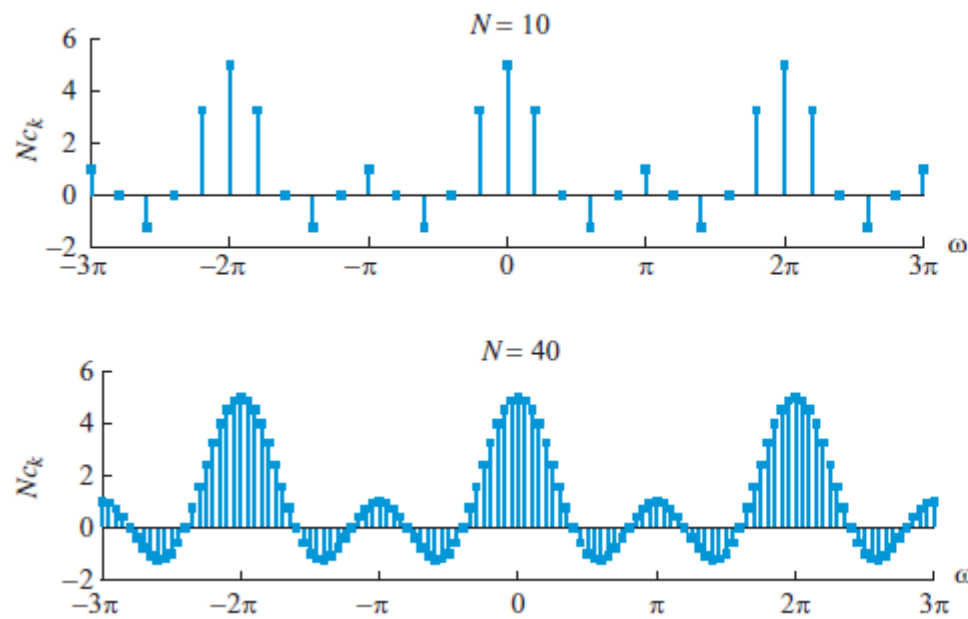
Fourier transforms for discrete-time aperiodic signals

- An **aperiodic** signal as a **periodic** signal with **infinite period**, we could obtain its Fourier representation by taking the limit of DTFS as the period increases indefinitely. For example let us start with the rectangular pulse train:

$$x[n] = \begin{cases} 1, & |n| \leq L \\ 0, & L < |n| < N/2 \end{cases} \quad \text{and periodically repeats with period } N > 2L + 1$$

$$c_k = \frac{1}{N} \frac{\sin\left[\frac{2\pi}{N} k \left(L + \frac{1}{2}\right)\right]}{\sin\left(\frac{2\pi}{N} k \frac{1}{2}\right)}$$

- As $N \rightarrow \infty$, $x[n]$ becomes an aperiodic sequence and its Fourier representation becomes a continuous function of ω .



- The discrete **Fourier transform representation (DTFT)** of **aperiodic signal**, is defines as follows:

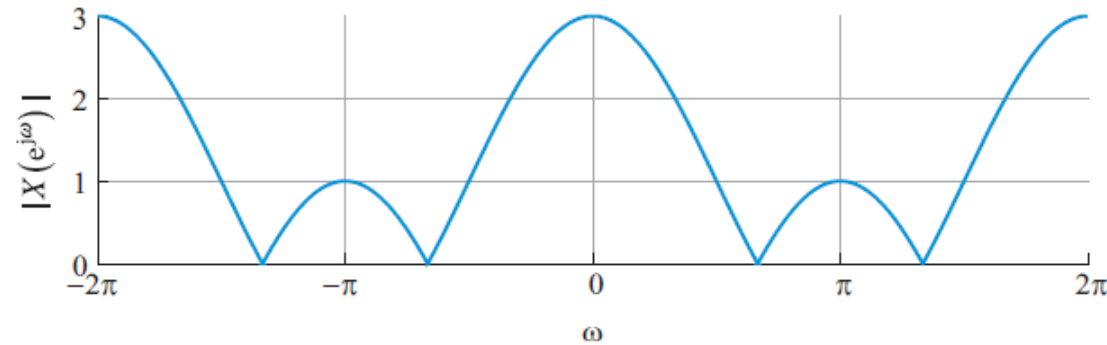
1. Synthesis equation: $x[n] = \mathcal{F}^{-1}\{\tilde{X}(e^{j\omega})\} = \frac{1}{2\pi} \int_{-\pi}^{\pi} \tilde{X}(e^{j\omega}) e^{j\omega n} d\omega$

2. Analysis equation: $\tilde{X}(e^{j\omega}) = \mathcal{F}\{x[n]\} = \sum_{n=-\infty}^{\infty} x[n] e^{-j\omega n}$

- **Example 1:** Finite length pulse

Evaluate and plot the magnitude and phase of the DTFT of the sequence $x[n] = \delta[n + 1] + \delta[n] + \delta[n - 1]$

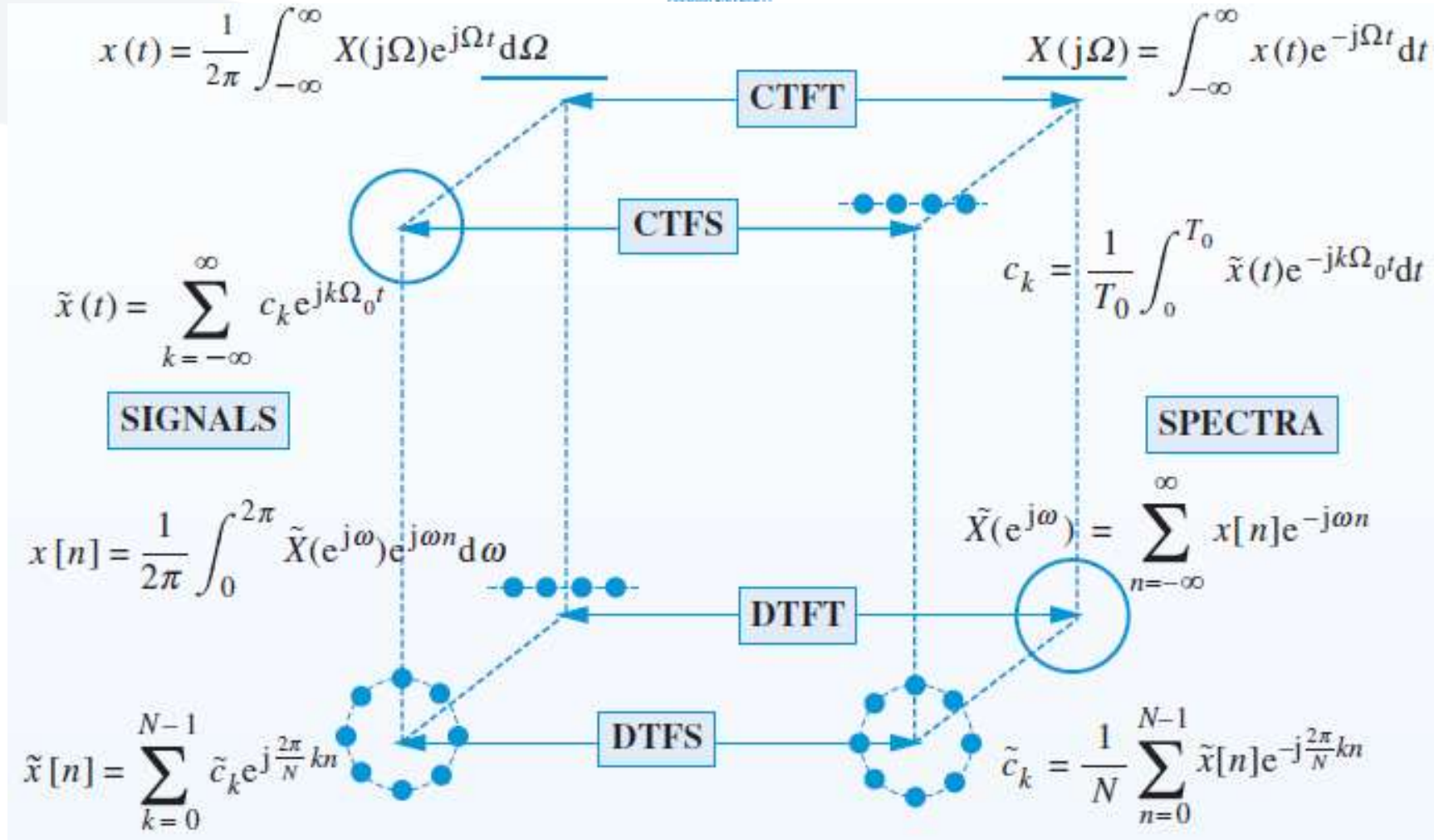
$$X(e^{j\omega}) = \sum_{n=-1}^1 x[n] e^{-j\omega n} = e^{j\omega} + 1 + e^{-j\omega} = 1 + 2 \cos(\omega)$$



- **Parseval's relation:** If $x[n]$ has **finite energy**, we have the following Parseval's relation:

$$E_x = \sum_{n=-\infty}^{\infty} |x[n]|^2 = \frac{1}{2\pi} \int_{2\pi} |\tilde{X}(e^{j\omega})|^2 d\omega$$

- The quantity $|\tilde{X}(e^{j\omega})/2\pi|^2$ or $|\tilde{X}(e^{j2\pi f})|^2$ is known as the **energy-density spectrum** of $x[n]$.
- There are four types of signal and related Fourier transform and series representations which are summarized in the figure below:



Summary of Fourier series and Fourier transforms

- **Continuous-time periodic signals** have **discrete aperiodic spectrum**. The spectrum exists only at $F = 0, \pm F_0, \pm 2F_0, \dots$, that is, at discrete values of F . The spacing between the lines of this discrete or line spectrum is $F_0 = 1/T_0$.
- **Continuous-time aperiodic signals** have **continuous aperiodic spectrum** over the entire frequency axis. The spectrum exists for all F , $-\infty < F < \infty$.
- **Discrete-time periodic signals** have **discrete periodic spectrum**. The spacing between the lines of the resulting discrete spectrum is $\omega = 2\pi/N$.
- **Discrete-time aperiodic signals** have **continuous periodic spectrum** of period 2π .
- **Note:** Periodicity with “period” α in one domain implies discretization with “spacing” of $1/\alpha$ in the other domain, and vice versa.

- **Bandlimited signals:** Signals whose frequency components are zero or “small” outside a finite interval $0 \leq B_1 \leq |F| \leq B_2 < \infty$ are said to be **bandlimited**. The quantity $B = B_2 - B_1$ is known as the **bandwidth** of the signal.
- **Bandlimited signals** cannot be **time-limited**, and vice versa; **Time-limited** signals cannot be **bandlimited**.

4. Properties of the discrete-time Fourier transform

Relationship to the z-transform and periodicity

- If the ROC of $X(z)$ includes the unit circle, defined by $z = e^{j\omega}$ or equivalently $|z| = 1$, we obtain:

$$X(z)\Big|_{z=e^{j\omega}} = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n} = \tilde{X}(e^{j\omega})$$

the z-transform reduces to the Fourier transform.

Symmetry properties

- Suppose that both the signal $x[n]$ and its DTFT $X(e^{j\omega})$ are complex-valued functions. Then: $x[n] = x_R[n] + jx_I[n]$, and $X(e^{j\omega}) = X_R(e^{j\omega}) + jX_I(e^{j\omega})$

$$\tilde{X}(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n] e^{-j\omega n} \quad \left\{ \begin{array}{l} X_R(e^{j\omega}) = \sum_{n=-\infty}^{\infty} \{x_R[n] \cos(\omega n) + x_I[n] \sin(\omega n)\} \\ X_I(e^{j\omega}) = - \sum_{n=-\infty}^{\infty} \{x_R[n] \sin(\omega n) - x_I[n] \cos(\omega n)\} \end{array} \right.$$

$$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} \tilde{X}(e^{j\omega}) e^{j\omega n} d\omega \quad \left\{ \begin{array}{l} x_R[n] = \frac{1}{2\pi} \int_{2\pi} [X_R(e^{j\omega}) \cos(\omega n) - X_I(e^{j\omega}) \sin(\omega n)] d\omega \\ x_I[n] = \frac{1}{2\pi} \int_{2\pi} [X_R(e^{j\omega}) \sin(\omega n) + X_I(e^{j\omega}) \cos(\omega n)] d\omega \end{array} \right.$$

- **Real signals** If $x[n]$ is real, then $x_R[n] = x[n]$ and $x_I[n] = 0$

$$X_R(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]\cos(\omega n) \text{ and } X_I(e^{j\omega}) = -\sum_{n=-\infty}^{\infty} x[n]\sin(\omega n)$$

$$\begin{aligned} X_R(e^{-j\omega}) &= X_R(e^{j\omega}) \quad (\text{even symmetry}) & X^*(e^{j\omega}) &= X(e^{-j\omega}) \\ X_I(e^{-j\omega}) &= -X_I(e^{j\omega}) \quad (\text{odd symmetry}) & & (\text{Hermitian symmetry}) \end{aligned} \Leftrightarrow$$

Thus, the DTFT of a real signal has **Hermitian** (or **complex-conjugate**) symmetry.

The inverse DTFT is given by: $x[n] = \frac{1}{2\pi} \int_{2\pi} [X_R(e^{j\omega})\cos(\omega n) - X_I(e^{j\omega})\sin(\omega n)] d\omega$

Since $X_R(e^{j\omega})\cos(\omega n)$ and $X_I(e^{j\omega})\sin(\omega n)$ are even functions of ω , we have

$$x[n] = \frac{1}{\pi} \int_0^{\pi} [X_R(e^{j\omega})\cos(\omega n) - X_I(e^{j\omega})\sin(\omega n)] d\omega$$

- **Real and even signals** If $x[n]$ is real and even, that is, $x[-n] = x[n]$, then:

$$\begin{cases} X_R(e^{j\omega}) = x[0] + 2\sum_{n=1}^{\infty} x[n]\cos(\omega n) & \text{(even symmetry)} \\ X_I(e^{j\omega}) = 0 \\ x[n] = \frac{1}{\pi} \int_0^{\pi} X_R(e^{j\omega})\cos(\omega n)d\omega & \text{(even symmetry)} \end{cases}$$

Thus, **real signals** with **even symmetry** have **real spectra** with **even symmetry**.

- **Real and odd signals** If $x[n]$ is real and odd, that is, $x[-n] = -x[n]$, then:

$$\begin{cases} X_R(e^{j\omega}) = 0 \\ X_I(e^{j\omega}) = -2\sum_{n=1}^{\infty} x[n]\sin(\omega n) & \text{(odd symmetry)} \\ x[n] = -\frac{1}{\pi} \int_0^{\pi} X_I(e^{j\omega})\sin(\omega n)d\omega & \text{(odd symmetry)} \end{cases}$$

Thus, **real signals** with **odd symmetry** have **purely imaginary** spectra with **odd symmetry**.

Property	$x[n]$	$X(e^{j\omega})$
Linearity	$ax_1[n] + bx_2[n]$	$aX_1(e^{j\omega}) + bX_2(e^{j\omega})$
Time shifting	$x[n - k]$	$X(e^{j\omega})e^{-jk\omega}$
Frequency shifting	$x[n]e^{j\omega_0 n}$	$X(e^{j(\omega-\omega_0)})$
Time reversal (Folding)	$x[-n]$	$X(e^{-j\omega})$
Conjugation	$x^*[n]$	$X^*(e^{-j\omega})$
Modulation	$x[n]\cos(\omega_0 n)$	$\frac{1}{2}[X(e^{j(\omega+\omega_0)}) + X(e^{j(\omega-\omega_0)})]$
Differentiation	$nx[n]$	$-j dX(e^{j\omega})/d\omega$
Convolution	$x_1[n] * x_2[n]$	$X_1(e^{j\omega}) X_2(e^{j\omega})$
Windowing	$x[n]w[n]$	$\frac{1}{2\pi} \int_{2\pi} X(e^{j\theta}) W[e^{j(\omega-\theta)}] d\theta$
Parseval's theorem	$\sum_{n=-\infty}^{\infty} x_1[n]x_2^*[n]$	$\frac{1}{2\pi} \int_{2\pi} X_1(e^{j\omega}) X_2^*(e^{j\omega}) d\omega$

- **Example 2:** Causal exponential sequence

Consider the sequence $x[n] = a^n u[n]$. For $|a| < 1$, the sequence is absolutely summable, that is:

$$\sum_{n=0}^{\infty} |a|^n = \frac{1}{1-|a|} < \infty$$

The DTFT exists and is given by: $X(e^{j\omega}) = \sum_{n=0}^{\infty} a^n e^{-j\omega n} = \frac{1}{1 - ae^{-j\omega}}$, $|a| < 1$

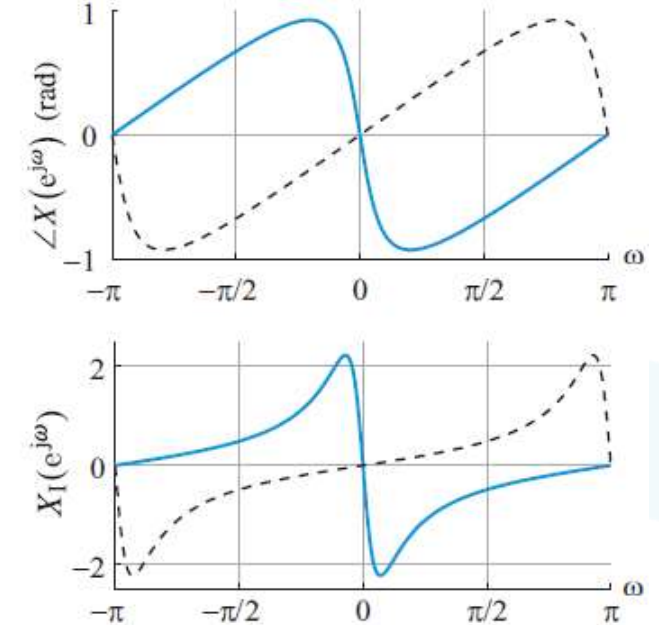
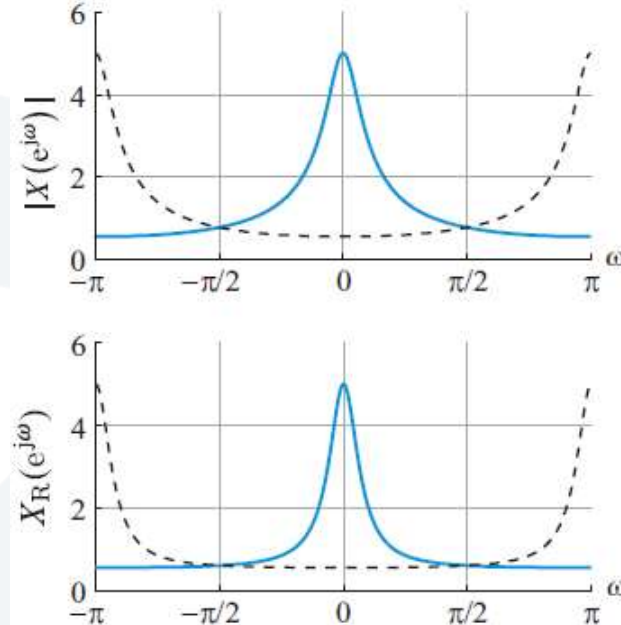
If $x[n]$ is real: $X_R(e^{j\omega}) = \frac{1 - a\cos(\omega)}{1 - 2a\cos(\omega) + a^2} = X_R(e^{-j\omega})$ (even)

$$X_I(e^{j\omega}) = \frac{-a\sin(\omega)}{1 - 2a\cos(\omega) + a^2} = -X_I(e^{-j\omega}) \quad (\text{odd})$$

$$|X(e^{j\omega})| = \frac{1}{\sqrt{1 - 2a\cos \omega + a^2}} = |X(e^{-j\omega})|, \quad \angle X(e^{j\omega}) = \tan^{-1} \frac{-a\sin \omega}{1 - a\cos \omega} = -\angle X(e^{-j\omega})$$

- **Note:** Lowpass sequence for $(0 < a < 1)$ and highpass sequence $(-1 < a < 0)$.

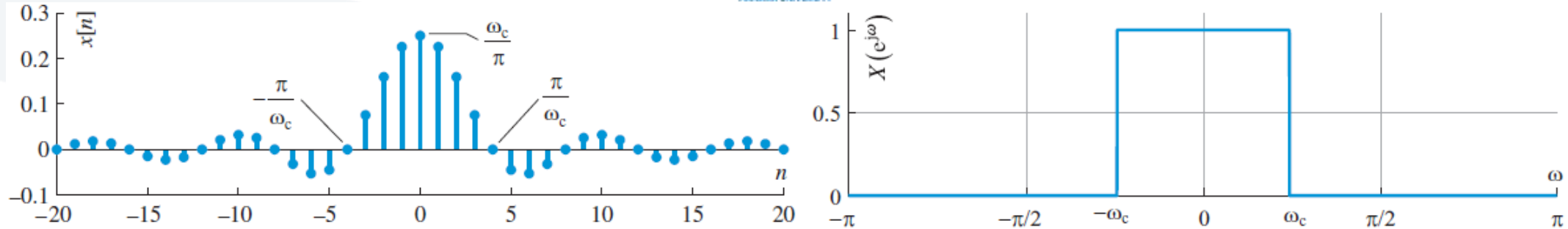
The solid lines correspond to a lowpass sequence ($a = 0.8$) and the dashed lines to a highpass sequence ($a = -0.8$)



- **Example 3:** Ideal lowpass sequence

A sequence $x[n]$ with DTFT over one period: $X(e^{j\omega}) = \begin{cases} 1, & |\omega| < \omega_c \\ 0, & \omega_c < \omega < \pi \end{cases}$

$$x[n] = \frac{1}{2\pi} \int_{-\omega_c}^{\omega_c} e^{j\omega n} d\omega = \frac{1}{2\pi jn} e^{j\omega n} \Big|_{-\omega_c}^{\omega_c} = \frac{\sin(\omega_c n)}{\pi n}, \quad -\infty < n < \infty$$



Correlation of signals

- To measure the **similarity** between a signal of **interest** and a **reference** signal, we use the **correlation sequence** of two real-valued signals, $x[n]$ and $y[n]$ each of which has finite energy, defined by:

$$\left. \begin{aligned} r_{xy}[l] &= \sum_{n=-\infty}^{\infty} x[n]y[n-l] = \sum_{n=-\infty}^{\infty} x[n+l]y[n], & -\infty < l < \infty \\ r_{yx}[l] &= \sum_{n=-\infty}^{\infty} y[n]x[n-l] = \sum_{n=-\infty}^{\infty} y[n+l]x[n], & -\infty < l < \infty \end{aligned} \right\} \Rightarrow r_{xy}[l] = r_{yx}[-l]$$

- Note:** We can compute correlation using convolution: $r_{xy}[l] = x[l] * y[-l]$

- To understand the meaning of correlation we first note that the energy E_z of the sequence $z[n] = ax[n] + y[n - l]$, which is nonnegative, can be expressed as:

$$E_z = a^2 E_x + 2ar_{xy}[l] + E_y \geq 0 \Rightarrow 4r_{xy}^2[l] - 4E_x E_y \leq 0$$

$$-1 \leq \rho_{xy}[l] = \frac{r_{xy}[l]}{\sqrt{E_x} \sqrt{E_y}} \leq 1$$

The sequence $\rho_{xy}[l]$, which is known as the **normalized correlation coefficient**:

- If $x[n] = cy[n - n_0]$, $c > 0$, we obtain $\rho_{xy}[n_0] = 1$ (maximum correlation);
- If $x[n] = -cy[n - n_0]$, $c > 0$, we obtain $\rho_{xy}[n_0] = -1$ (maximum negative correlation).
- If $\rho_{xy}[l] = 0$ for all lags, the two sequences are said to be **uncorrelated**.

$$r_{xy}[l] = x[l] * y[-l] \Rightarrow \text{using DTFT } R_{xy}(\omega) = X(e^{j\omega}) Y(e^{-j\omega})$$

when $y[n] = x[n]$ we obtain the **autocorrelation** sequence $r_{xx}[l]$ or $r_x[l]$.

Since $x[n]$ is a real sequence, $X^*(e^{j\omega}) = X(e^{-j\omega})$ therefore the DTFT of $r_x[l]$ is:

$$r_x[l] = x[l] * x[-l] \Rightarrow R_x(\omega) = |X(e^{j\omega})|^2 \quad \text{Wiener-Khintchine theorem}$$

- **Example 4:** Autocorrelation of exponential sequence

Let $x[n] = a^n u[n]$, $-1 < a < 1$. For $l > 0$, the product $x[n]u[n]x[n-l]u[n-l]$ is zero for $n < l$.

$$r_x[l] = \sum_{n=l}^{\infty} x[n]x[n-l] = \sum_{n=l}^{\infty} a^n a^{n-l} = a^l (1 + a^2 + a^4 + \dots) = \frac{a^l}{1-a^2}$$

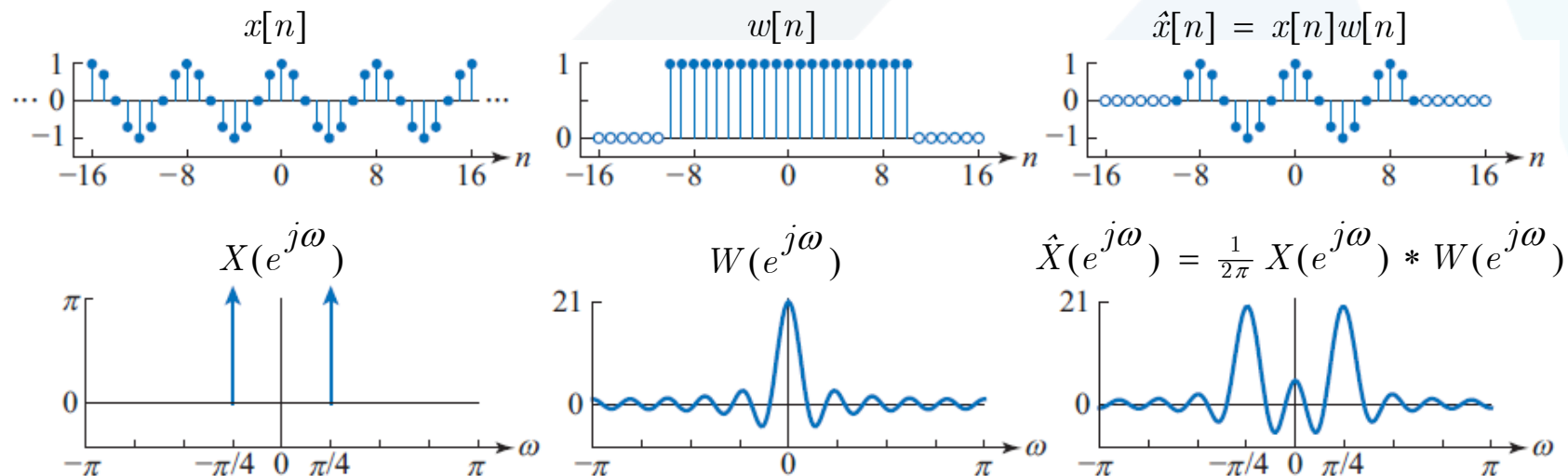
Since $r_x[l] = r_x[-l] \Rightarrow r_x[l] = \frac{a^{|l|}}{1-a^2}$, $-1 < a < 1$ The Fourier transform is

$$R_x(\omega) = X(e^{j\omega})X(e^{-j\omega}) = \frac{1}{1-ae^{-j\omega}} \frac{1}{1-ae^{j\omega}} = \frac{1}{1-2a \cos(\omega) + a^2}$$

Since $r_x[l]$ is real and even, its Fourier transform $R_x(\omega)$ is also real and even.

Spectral and temporal ambiguity

- **Spectral analysis** is one of the most important applications of DSP. It is the process of **measuring**, **estimating** and **characterizing** the **frequency** content of signals.
- Let $x[n] = \cos \pi n/4$, which might represent the signal whose spectrum we are trying to measure.



- $x[n]$ is completely **unlocalized** in time (extends from $-\infty < n < \infty$); however, its spectrum $X(e^{j\omega})$ is highly **localized** to exactly two frequencies, $\omega = \pm\pi/4$.
- The result of this **convolution** (windowing of a cosine) is a signal whose spectrum is “**smear**ed out” in frequency; it has two relatively broad peaks and **energy spread** over all frequencies.

- There is a **trade-off** between resolution in the time and the frequency domains; as we **increase** the **localization** of the signal in the **time domain**, we **reduce** the localization in the frequency domain.

